

New tools for conformal manifolds in one and two dimensions

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Conformal field theory

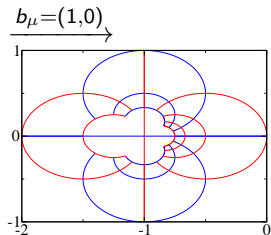
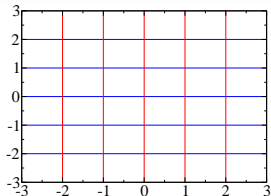
More symmetric than a “typical” QFT:

Translations	$x'_\mu = x_\mu + a_\mu$
Rotations	$x'_\mu = \Lambda_\mu^\nu x_\nu$
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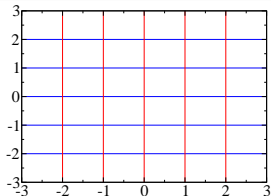
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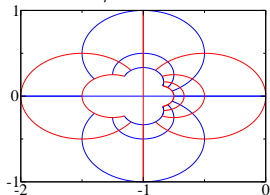
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$$\mathcal{O}_{\mu_1 \dots \mu_\ell}(x) \mapsto \lambda^{-\Delta} \mathcal{O}_{\mu_1 \dots \mu_\ell}(\lambda x)$$



$b_\mu = (1, 0)$
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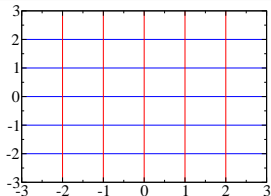
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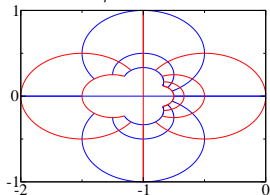
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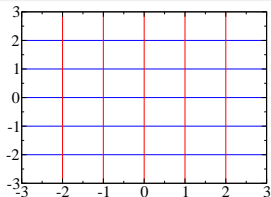
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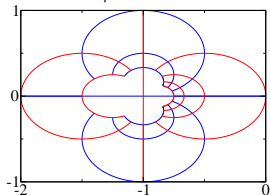
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- $\langle \phi(x_1) \phi(x_2) \rangle = \frac{1}{|x_{12}|^{2\Delta}}$
- $\langle \phi_1(x_1) \phi_2(x_2) \phi_3(x_3) \rangle = \frac{\lambda_{123}}{|x_{12}|^{\Delta_1 + \Delta_2 - \Delta_3} |x_{23}|^{\Delta_2 + \Delta_3 - \Delta_1} |x_{13}|^{\Delta_1 + \Delta_3 - \Delta_2}}$
- $\phi_1(x_1) \phi_2(x_2) = \sum_{\mathcal{O}} \frac{\lambda_{12\mathcal{O}}}{|x_{12}|^{\Delta_1 + \Delta_2 - \Delta}} C_{(\mu)}(x_{12}, \partial_2) \mathcal{O}^{(\mu)}(x_2)$



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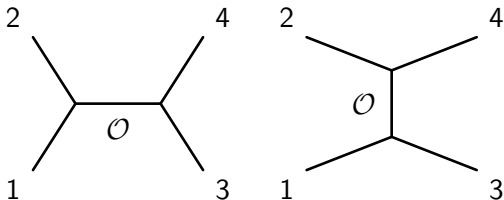


Conformal bootstrap

$\langle \phi_1(x_1)\phi_2(x_2)\phi_3(x_3)\phi_4(x_4) \rangle$ includes unknown function of
 $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}$ and $v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$ but there are two ways to compute it!

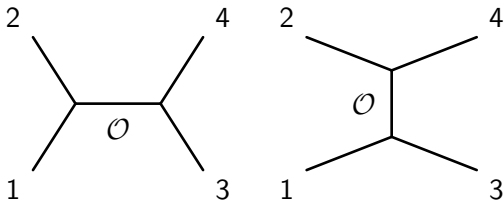
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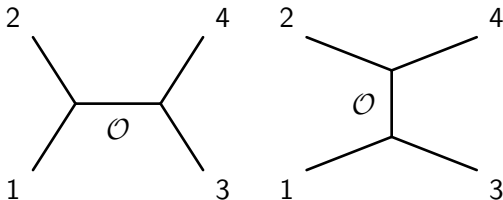
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Sum over \mathcal{O} must include operators in certain ranges for crossing symmetry to have a solution. [\[Rattazzi, Rychkov, Tonni, Vichi; 0807.0004\]](#)

$$\sum_{\mathcal{O}} \lambda_{\phi\phi\mathcal{O}}^2 [v^{\Delta_{\phi}} G_{\mathcal{O}}(u, v) - u^{\Delta_{\phi}} G_{\mathcal{O}}(v, u)] = 0$$

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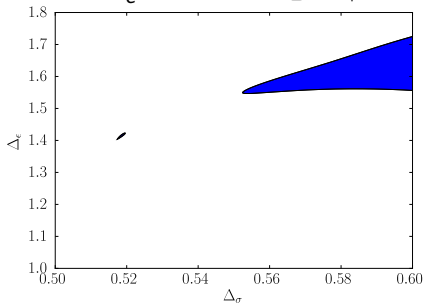
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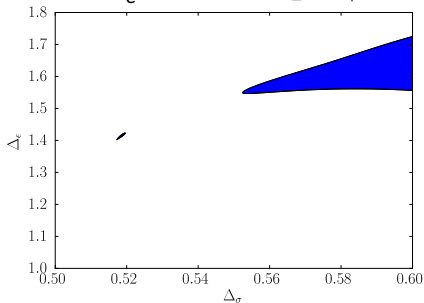
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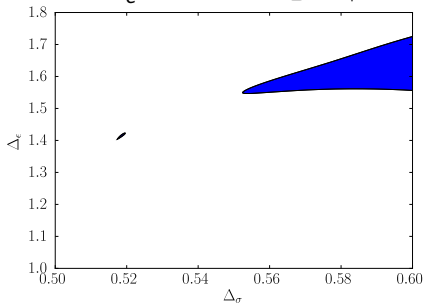
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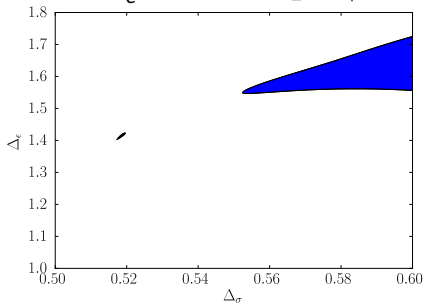
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Perturbative control for large and small s . [CB, Rastelli, Rychkov, Zan; 1703.05325]

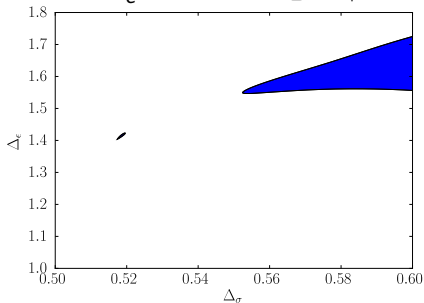
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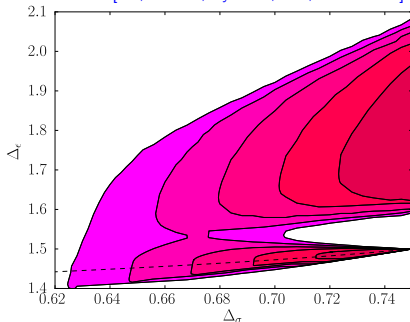


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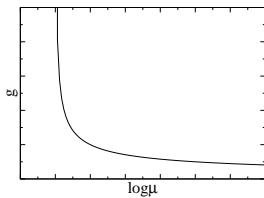
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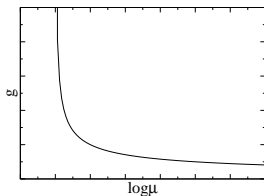
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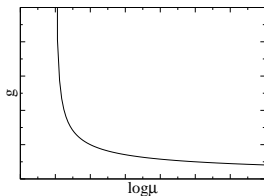
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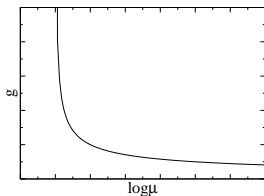
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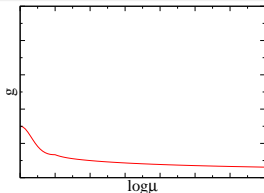
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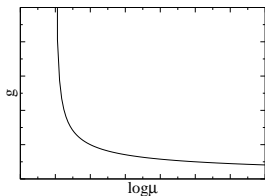
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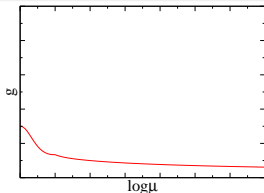
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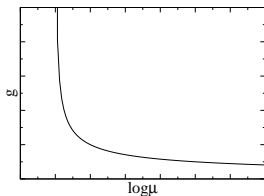


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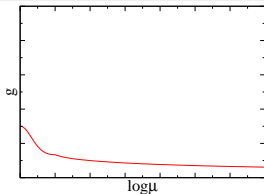
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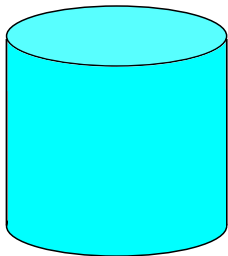
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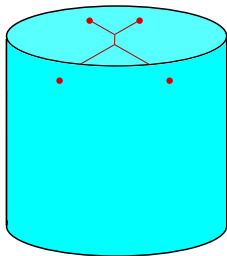
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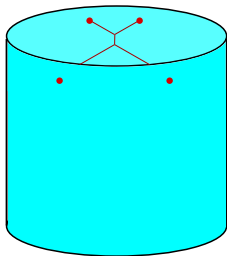
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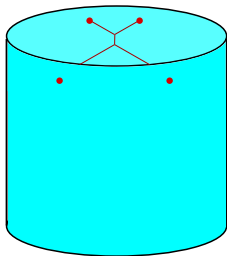
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$\Delta = d$ from symmetry breaking.

[Carmi, Di Pietro, Komatsu; 1810.04185]



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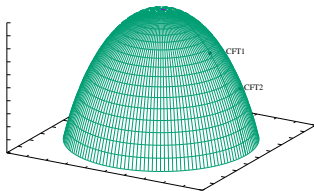
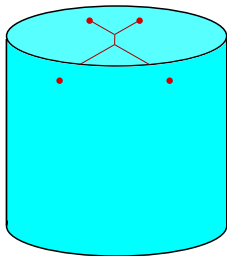
$\mathcal{N} = 4$ Super Yang-Mills	λ $\tilde{\lambda}$	A_μ	φ $\tilde{\varphi}^\dagger$	ψ
---------------------------------------	--------------------------------	---------	--	--------

$$ds^2 = \frac{L^2}{z^2} (dz^2 + \eta_{\mu\nu} dx^\mu dx^\nu)$$

$$(mL)^2 = \Delta(\Delta - d)$$

$\Delta = d$ from symmetry breaking.

[Carmi, Di Pietro, Komatsu; 1810.04185]



Flow by adding $\int \delta g \hat{O} dx$

$$\delta \Delta_i = \delta g A_i(\{\Delta\}, \{\lambda\}) + O(\delta g^2)$$

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$$\langle \mathcal{O}_1 \dots \mathcal{O}_n \rangle_{g+\delta g} = \left\langle \mathcal{O}_1 \dots \mathcal{O}_n \exp \left[\int \delta g \hat{\mathcal{O}} dx \right] \right\rangle_g$$

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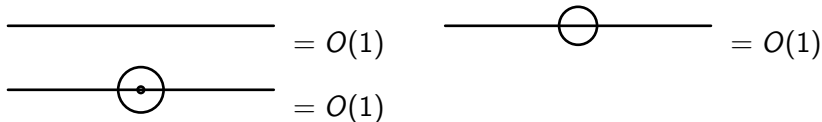
$$S = \int \frac{1}{2} \chi_i \frac{d}{d\tau} \chi^i + \frac{1}{4!} J_{ijkl} \chi^i \chi^j \chi^k \chi^l d\tau, \quad \langle J_{ijkl} J^{ijkl} \rangle = 3! J^2 / N^3$$

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$$\text{—————} = O(1) \quad \text{—————} \bigcirc \text{—————} = O(1)$$

SYK-like models

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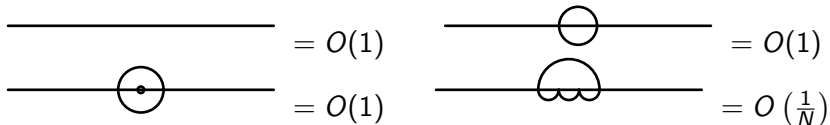


The image shows four Feynman diagrams arranged in two rows. The top row consists of two diagrams: a horizontal line on the left, and a horizontal line on the right with a circle in the middle. The bottom row consists of two diagrams: a horizontal line on the left with a circle in the middle containing a central dot, and a horizontal line on the right with a circle in the middle. Each diagram is followed by an equals sign and the expression $O(1)$.

$$\begin{array}{ll} \text{---} & = O(1) \\ \text{---} & = O(1) \end{array}$$

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$\text{Diagram 1} = O(1)$ $\text{Diagram 2} = O(1)$
 $\text{Diagram 3} = O(1)$ $\text{Diagram 4} = O\left(\frac{1}{N}\right)$

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 \text{---} \circ \text{---} & = O(1) \\
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 \end{array}$$

Kinetic term $\mapsto \int \text{sgn}(\tau_{12}) \chi_i(\tau_1) \chi^i(\tau_2) \frac{d\tau_1 d\tau_2}{|\tau_{12}|^{\frac{3}{2}}}$. [\[Gross, Rosenhaus; 1706.07015\]](#)

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$$\beta_0 \sim \frac{J_1 J_2}{N^{\frac{5}{2}}}, \quad \beta_1 \sim \frac{J_0 J_3}{N^{\frac{5}{2}}}, \quad \beta_2 \sim \frac{J_1^2 + J_2^2 + J_3^2}{N^2}, \quad \beta_3 \sim \frac{J_0^2 + J_3^2 + J_2 J_3}{N^2}$$

Free boson with a boundary deformation

Pieces of $\Phi(z, \bar{z}) = X(z) + \bar{X}(\bar{z})$ agree on Neumann boundary.

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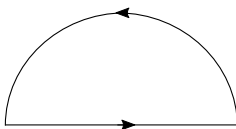
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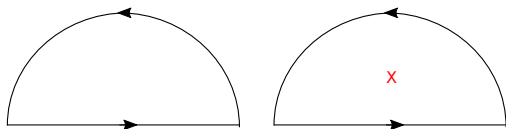
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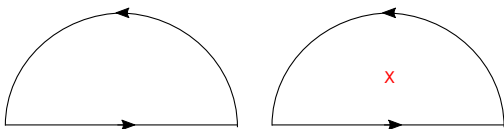
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One-point $\langle \mathcal{O} \rangle \equiv \langle \mathcal{O} | B \rangle$ can be combined into boundary state.

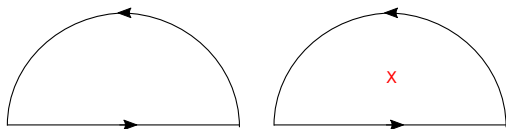
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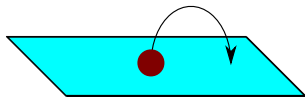


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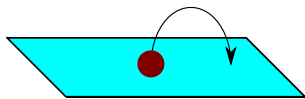
$$Z_{[0,L] \times [0,T]} = \left\langle B \left| e^{-\frac{2\pi T}{L} H} \right| B \right\rangle$$

[Callan, Klebanov, Ludwig, Maldacena; hep-th/9402113]

Instantons (with Rastelli, Roumpedakis)

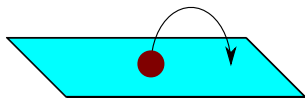


Instantons (with Rastelli, Roumpedakis)



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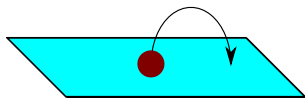


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Dimension-1 operator producing twisted states:

$$V_w + V_{\bar{w}} = w_\alpha G_{-\frac{1}{2}} \Delta S^\alpha(z) + \bar{w}_\alpha G_{-\frac{1}{2}} \bar{\Delta} S^\alpha(z)$$

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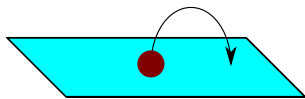
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Should control moduli in instanton solution:

$$A_\mu(x)^\alpha_\beta = (\bar{\sigma}_{\mu\nu})^\alpha_\beta \frac{\rho^2 x^\nu}{x^2(x^2 + \rho^2)}$$

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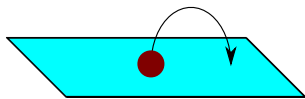
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$$A_\mu(x)_\beta^\alpha = (\bar{\sigma}_{\mu\nu})_\beta^\alpha \frac{\rho^2 x^\nu}{x^2(x^2 + \rho^2)} = (\bar{\sigma}_{\mu\nu})_\beta^\alpha \frac{\rho^2 x^\nu}{x^4} + \dots$$

[Billo, Frau, Fucito, Lerda, Liccardo, Pesando; hep-th/0211250]

Instantons (with Rastelli, Roumpedakis)



- $X^\mu(z)$ anti-periodic around $\Delta(z)$, $\bar{\Delta}(z)$
- $\psi^\mu(z)$ anti-periodic around $S^\alpha(z)$

Dimension-1 operator producing twisted states:

$$V_w + V_{\bar{w}} = w_\alpha G_{-\frac{1}{2}} \Delta S^\alpha(z) + \bar{w}_\alpha G_{-\frac{1}{2}} \bar{\Delta} S^\alpha(z)$$

Should control moduli in instanton solution:

$$A_\mu(x)_\beta^\alpha = (\bar{\sigma}_{\mu\nu})_\beta^\alpha \frac{\rho^2 x^\nu}{x^2(x^2 + \rho^2)} = (\bar{\sigma}_{\mu\nu})_\beta^\alpha \frac{\rho^2 x^\nu}{x^4} + \dots$$

[Billo, Frau, Fucito, Lerda, Liccardo, Pesando; hep-th/0211250]

Root vectors of $SO(8)$ appear at $R = \sqrt{2}$!

$$V_w + V_{\bar{w}} = \exp [i\alpha \cdot (\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6)]$$

- Several theories can benefit from more bootstrap-like approaches.
- A system of differential equations controls the data along a conformal manifold.
- Promising targets include SYK-like models, BCFTs in $d = 1 + 1$ and the D-instanton.
- Non-linear sigma models, $\mathcal{N} = 4$ SYM and AdS Goldstones could be next.